

3.5. Feynman theory

Consider a smooth 2D vector field $V(\varphi)$. The angle that vector V makes with respect to $\hat{\varphi}_1$ and $\hat{\varphi}_2$ axes is a scalar field:

$$\theta(\varphi) = \tan^{-1} \left(\frac{V_2(\varphi)}{V_1(\varphi)} \right)$$

As long as V has finite length [doesn't vanish or blow up] the angle θ is well defined. Expect that we can integrate $\nabla \theta$ over a closed curve C in phase space to get

$$\oint_C d\varphi \cdot \nabla \theta = 0$$

[θ at all points is independent of path taken]

[$\nabla \theta$ is gradient of a scalar field]

However this can fail if $V(\varphi)$ vanishes or diverges at one or more pts. in the interior of C .

Define:
$$W_C(V) = \frac{1}{2\pi} \oint_C d\varphi \cdot \nabla \theta$$

then $W_C(V) \in \mathbb{Z}$ is an integer valued function of C which is the change in θ around the curve C . This must be an integer because θ is well defined only upto multiples of 2π . Differential changes of θ are in general well defined.

If $V(\varphi)$ is finite \rightarrow neither diverges nor vanishes anywhere in $\text{int}(C)$ then $W_C(V) = 0$.

Assuming that V never diverges, any singularities in θ must be arising from pts where $V=0$, which in general occurs at isolated points. since two variables

The index of a two-dimensional vector field $V(\varphi)$ at a point φ is the integer-valued winding of V about that point.

$$\text{ind}(V)_\varphi = \lim_{a \rightarrow 0} \frac{1}{2\pi} \oint_{C_a(\varphi_0)} d\varphi \cdot \nabla \theta$$

$$= \lim_{a \rightarrow 0} \frac{1}{2\pi} \oint d\varphi \cdot \left(\frac{V_1 \nabla V_2 - V_2 \nabla V_1}{V_1^2 + V_2^2} \right)$$

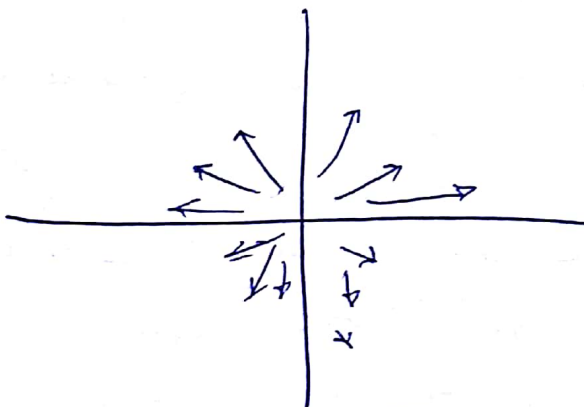
The index of a closed curve C is given by the sum of the indices at all the singularities enclosed by the curve*.

$$\left[\because \theta = \tan^{-1} \left(\frac{V_2}{V_1} \right) \right]$$

$$d\theta = \frac{1}{1 + \left(\frac{V_2}{V_1} \right)^2} \times \left(\frac{V_1 dV_2 - V_2 dV_1}{V_1^2} \right)$$

$$W_C(V) = \sum_{\varphi_i \in \text{int}(C)} \text{ind}(V)_{\varphi_i}$$

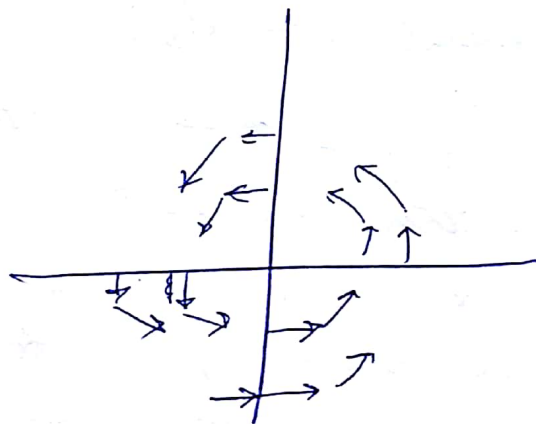
* Technically should weight the index of each enclosed singularity by the signed number of times the curve C encloses that singularity. For simplicity and clarity, assume that curve C is homeomorphic to the circle S^1 .



Index

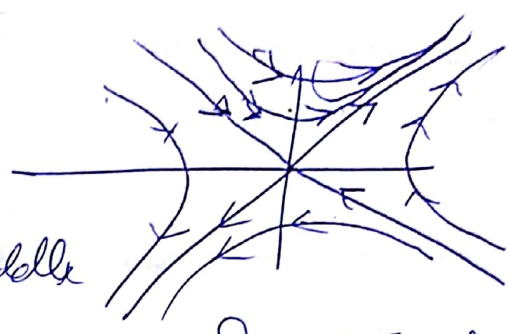
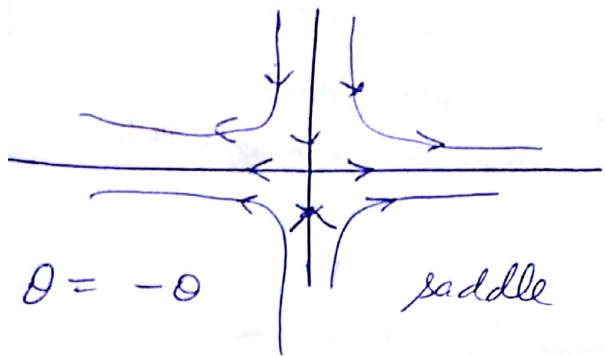
+1

Unstable Node.



+1

Periodic cycle
(Purely Imaginary)



$\theta = -0$

saddle

saddle

$\theta = \frac{\pi}{2} - 0$

Index -1

-1

1st diagram $v = (x, y) \rightarrow \theta = 0$

$v = (-y, x) \rightarrow \theta = 0 + \frac{\pi}{2}$

+1

2nd $v = (x, -y)$
 $v = (y, x)$

$\rightarrow \theta = -0$

$\rightarrow \theta = -0 + \frac{\pi}{2}$

-1

$v = (x^2 - y^2, 2xy)$

$v = (1 + x^2 - y^2, x + 2xy)$

$\theta = 2\theta$

+2

single fixed pt at $(0,0)$ index +2

two fixed pts:

$(x, y) = \begin{cases} (0, 1) \\ (0, -1) \end{cases}$

$M = \begin{pmatrix} 2x & -2y \\ 1+2y & 2x \end{pmatrix}$

$M_0 = \begin{pmatrix} 0 & -2 \\ 2 & 0 \end{pmatrix}$ $M_{0-1} = \begin{pmatrix} 0 & 2 \\ 2 & 0 \end{pmatrix}$

each +1

+1

$\therefore w_c(v) = +1 + (+1) = +2$

Properties of index/winding no.

- The index ν of $N=2$ vector field v at pt φ_0 is winding no. of v about that point.

- The winding number $w_c(v)$ of a curve c is the sum of indices of the singularities enclosed by that curve.

- Smooth deformations of c do not change the winding no. One must stretch c over a fixed pt singularity in order to change $w_c(v)$.

- Uniformly rotating each vector in the vector field by angle β has the effect of sending $\theta \rightarrow \theta + \beta$. This leaves all indices and winding no.s invariant.

- Nodes and spirals whether stable or unstable have index of +1 (as do the special cases of centers, stars and degenerate nodes)
- Saddle points have index -1.

- Clearly any closed orbit must lie on a curve C of index +1.

3.6.1 Gauss-Bonnet Theorem

Deep result in mathematics: connects local geometry of 2D manifold to global topological structure.

Content of theorem:

$$\int_M dA K = 2\pi \chi(M) = 2\pi \sum_i \text{ind}(\psi_i)$$

where M is a 2D manifold (a topological space locally homeomorphic to \mathbb{R}^2)
 K is the local gaussian curvature of M :

$$K = \frac{1}{R_1 R_2}$$

where R_1, R_2 are principal radii of curvature.

$\chi(M)$: Euler characteristic of M

given by $\chi(M) = 2 - 2g$ where g is the genus of M , which is the no. of holes (or handles) of M .

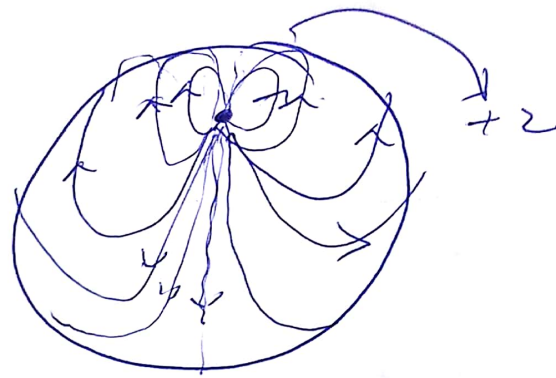
ψ is any smooth vector field on M , ψ_i are the singularity points of the vector field, which are fixed pts of the dynamics $\dot{\psi} = \psi(\psi)$

Consider $M = S^2$ unit 2-sphere.

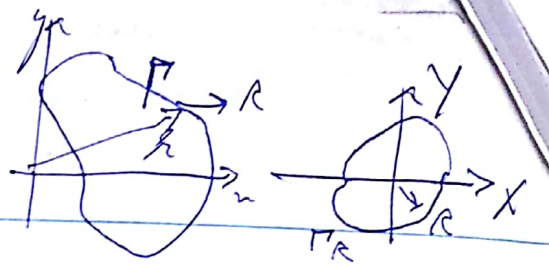
at any pt $R_1 = R_2 = 1 \rightarrow K = 1$ $\int dA K = 4\pi$

$$4\pi = (2\pi) \chi(M) \rightarrow \chi(M) = 2$$

Also any smooth vector field on S^2 must have a singularity or singularities. $\sum \text{ind}(\psi) = \chi(S^2) = 2$



$$\frac{d\theta}{ds} = \frac{x\dot{y} - y\dot{x}}{x^2 + y^2}$$



$$I\Gamma = \frac{1}{2\pi} \oint \frac{d\theta}{ds} ds = \frac{1}{2\pi} \oint \frac{x\dot{y} - y\dot{x}}{x^2 + y^2} ds$$

$$I\Gamma = \frac{1}{2\pi} \oint x dy - y dx$$

$r(s) = (x(s), y(s))$ describes Γ

$R(s) = (x, y)$ can be regarded as a position vector on a plane with x, y axis describes another curve Γ_R

Γ_R is closed since R returns to its original value after a complete cycle

Γ_R encircles origin $I\Gamma$ times, anticlockwise

Thm 3.1 Suppose on and inside $\Gamma: x, y$ and their 1st derivatives are cont and x and y are not simultaneously 0. (No eq/pt) then $I\Gamma$ is 0.

Green's Thm

$$I\Gamma = \frac{1}{2\pi} \oint \left[\frac{\partial}{\partial x} \left(\frac{y}{x^2 + y^2} \right) - \frac{\partial}{\partial y} \left(\frac{x}{x^2 + y^2} \right) \right] dx dy$$

\int_{SR} region enclosed by Γ_R

Integrand is identically 0.

($u_x = -v_y$)

Green's theorem

$$\oint_C P dx + Q dy = \iint_D \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dA$$

Index Theory

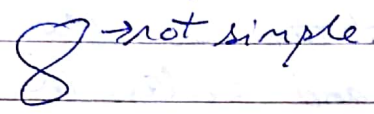
Strogatz NLI
lec 8 Youtube

Provides "global" info about phase portrait.

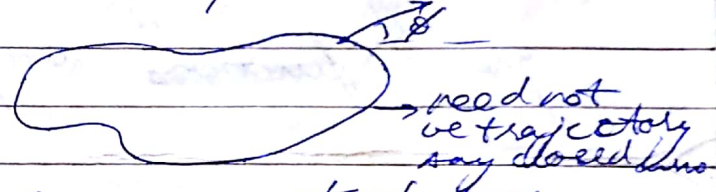
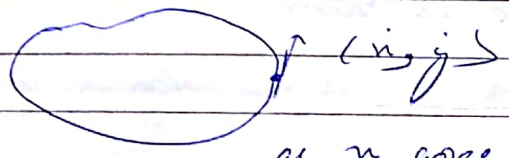
Linearizing: local methods \rightarrow However depends on info in ^{extremal} neighborhood of pt.

Topology: complex - winding No. Physics: Gauss.

Index of a closed curve c \rightarrow (doesn't intersect itself) ^{and doesn't pass through fixed pt.}
 $c =$ simple closed curve, not necessarily a closed trajectory.



$$\phi = \tan^{-1} [y/x]$$



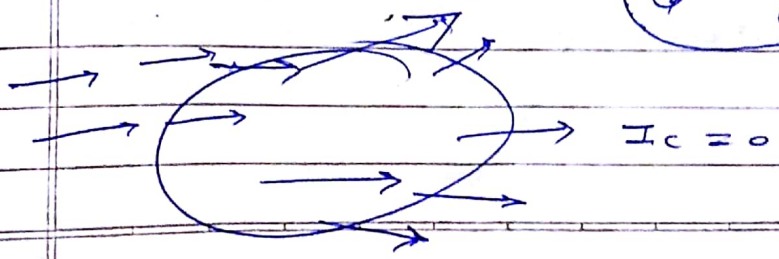
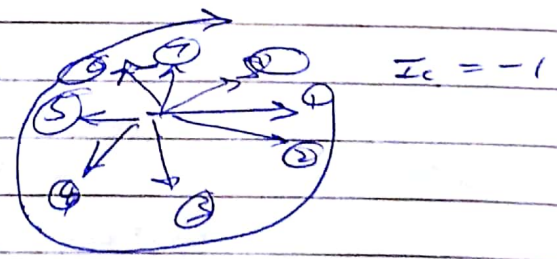
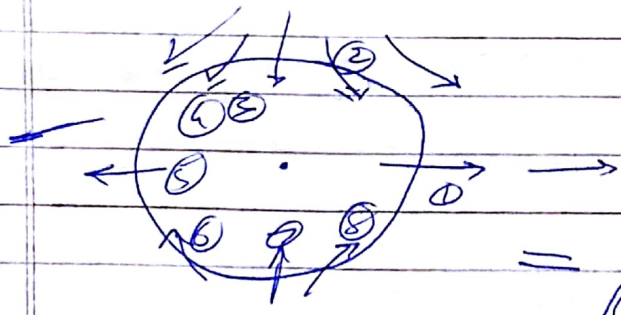
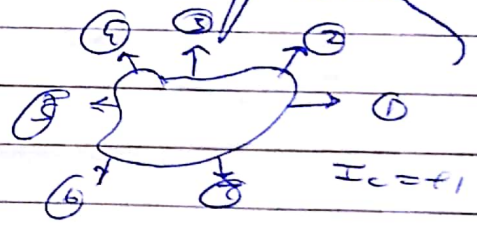
as n goes around c once, counterclockwise,

ϕ changes continuously.

If x, y are continuous, let $\sum \phi|_c =$ net change in ϕ around c . Then $I_c = \frac{1}{2\pi} \sum \phi|_c$ is index of c w.r.t vector field (x, y) .

$$\sum \phi|_c = 2\pi$$

$$I_c = +1$$

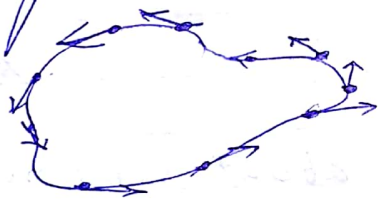


Teacher's Sign.: _____

stable or unstable

Properties:

↳ Index of a closed trajectory



Any closed trajectory has index $I_c = +1$

Questions?

Berry phase curvature and relⁿ to index theory



Brouwer's Fixed Point theorem and Sperner's lemma } Topology
triangulation and index theory

Green's functions, divergence, Stokes' theorem

Properties of the Index

1. Suppose C can be deformed continuously into C' without passing through a fixed pt. Then $I_C = I_{C'}$

(Deforming C into C' continuously then $\int_C \frac{d\phi}{2\pi} = \int_{C'} \frac{d\phi}{2\pi} =$ changes continuously. But I_C is an integer hence it can't change without jumping

2. If C does not enclose any fixed points, then $I_C = 0$

By property 1, shrink C to a tiny circle without changing index. But ϕ is essentially constant on such a circle, because all vectors point in nearly same direction thanks to the assumed smoothness of the vector field. Hence $\int_C \frac{d\phi}{2\pi} = 0$ $I_C = 0$

3. If we reverse all arrows in vector field by changing $t \rightarrow -t$, the index is unchanged.

$$\dot{x} = \frac{dx}{dt} = f(x) \quad \text{let } dt = -dt' \rightarrow \frac{dx}{dt'} = -f(x)$$

all angles change from ϕ to $\pi + \phi$ $\int_C \frac{d\phi}{2\pi}$ stays

4. Suppose closed curve C is actually a trajectory for system, i.e. C is a closed orbit. Then $I_C = +1$

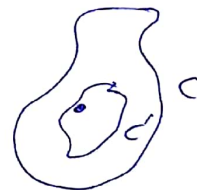
vector field is everywhere tangent to C because C is a trajectory



Index of a pt.

Let x^* be an isolated fixed pt.

I_C is independent of C hence only dependent on x^* .



* Index not related to stability

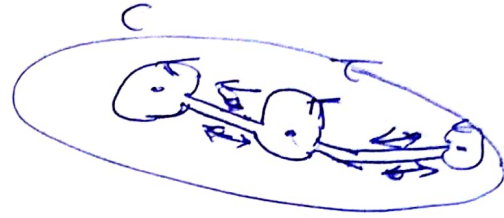
Spirals, centers, degenerate nodes and stars all have $I = +1$

* saddle pt. is truly different.

Thm: If a closed curve C surrounds n isolated
fixed pts. $n_1^*, n_2^*, \dots, n_n^*$. then

$$I_C = I_1 + I_2 + \dots + I_n$$

I_k is index of n_k^* for $k = 1, \dots, n$



Thm Any closed orbit in the phase plane must enclose
fixed pts whose indices sum to $+1$.

Corollary: Always at least one fixed pt inside
any closed orbit in the phase plane.

If there is only one fixed point inside it, it cannot
be a saddle point.

In every case the change in ϕ must be a multiple of 2π .

$$\oint_{\Gamma} \phi = 2\pi I_{\Gamma}$$

3.4

where I_{Γ} is defined and is an integer.

Index of Γ with ϕ vector field (x, y)

(Γ is always counterclockwise)

* Algebraic representation of I_{Γ}

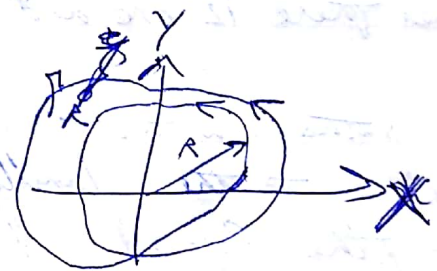
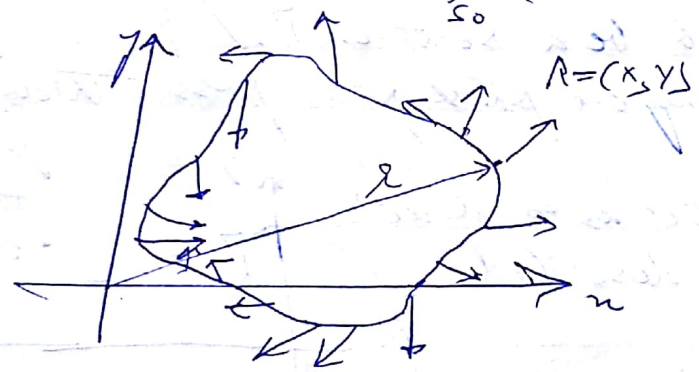
$$R(s) = (x(s), y(s)) \quad s_0 < s < s_1$$

$$\frac{d \tan \phi}{ds} = \frac{d}{ds} \left(\frac{y}{x} \right)$$

$$\therefore \frac{d\phi}{ds} = \frac{xy' - yx'}{x^2 + y^2}$$

from 3.4

$$I_{\Gamma} = \frac{1}{2\pi} \int_{s_0}^{s_1} \left(\frac{d\phi}{ds} \right) ds = \frac{1}{2\pi} \int_{s_0}^{s_1} \frac{xy' - yx'}{x^2 + y^2} ds$$



As $R(s)$ describes Γ , $R(s) = (x, y)$ describes another curve Γ_R on x, y axes plane.

Γ_R is closed since it returns to original value after complete cycle.

Γ_R encircles origin I_{Γ} times.

$$I_{\Gamma} = \frac{1}{2\pi} \oint_{\Gamma_R} \frac{x dy - y dx}{x^2 + y^2}$$

$$\left(\frac{dy}{dt} - \frac{dx}{dt} \right) dt$$

$$\left(dy - dx \right) \frac{dt}{dt}$$

take parameter $s = t$
 $ds = dt$

Thm Suppose Γ lies in a simply connected region on which x, y and their 1st derivatives are continuous and x and y are not simultaneously 0 (No origin).

then I_{Γ} is 0

Proof: Green's thm

$$\oint P du + Q dy = \iint_{\Delta \Gamma} \left(\frac{\partial Q}{\partial u} - \frac{\partial P}{\partial y} \right) du dy$$

We can write $dx = x_u du + x_y dy$ same for y
then $I_{\Gamma} = \frac{1}{2\pi} \oint \left(\frac{xy_u - yx_u}{x^2 + y^2} du + \frac{xy_y - yx_y}{x^2 + y^2} dy \right)$

Chapter 3 Geometrical Aspects of Plane Autonomous systems

Intro where the linear approximation is 0, Jordan, Smith index of a point provides supporting informⁿ on its nature and complexity in strongly nonlinear cases.

Phase diagram doesn't give info about behaviour of paths at infinity beyond boundaries (Global path S).

3.1 The index of a point.

Given the system: $\dot{x} = X(x, y)$ $\dot{y} = Y(x, y)$ (3.1)

let Γ be any smooth closed curve consisting of ordinary points of (3.1). let Q be a point on Γ .

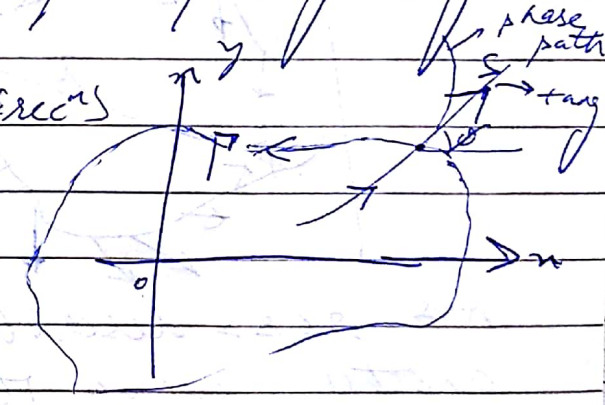
Then there is one and only one phase path passing through Q .

The paths (without implicit as to directions) belong to the family described by the eqn:

$$\frac{dy}{dx} = \frac{Y(x, y)}{X(x, y)}$$

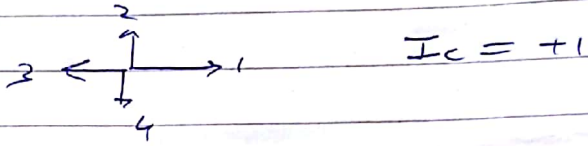
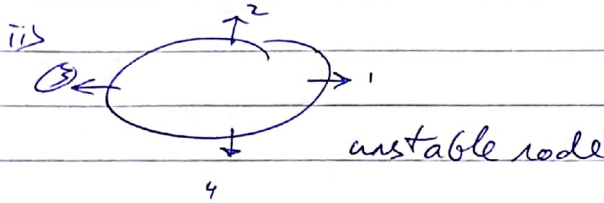
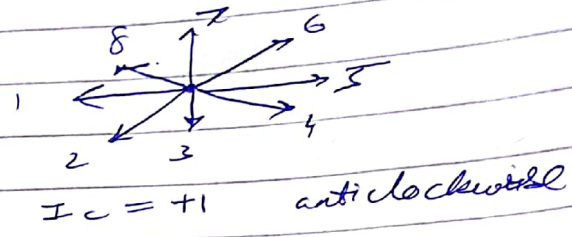
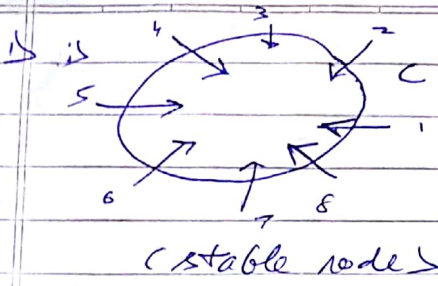
$S = (x, y)$ is tangential to phase path through the point and points in direction of increasing t .

$\tan \phi = Y/X$

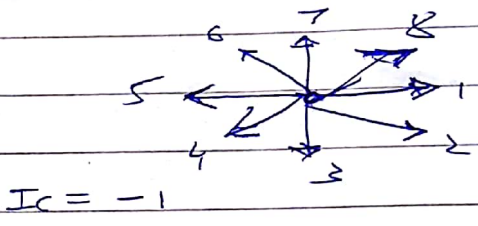
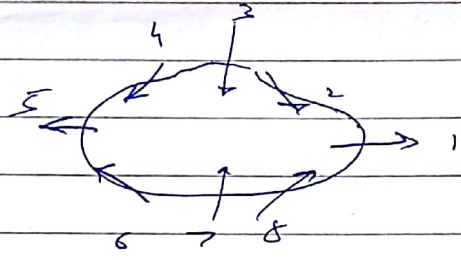


The curve Γ is traversed in the counterclockwise direction and variation of ϕ is followed along it ϕ will not return to original value but differ by $2\pi n$ after a full cycle.
Example.

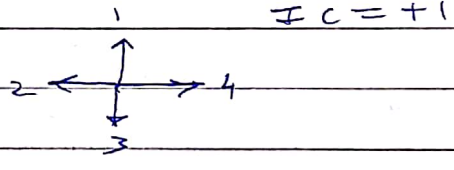
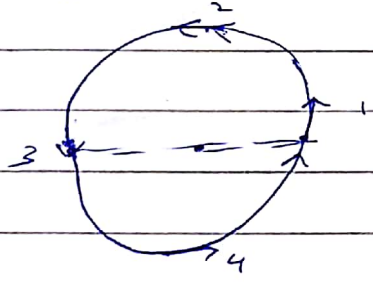
Examples



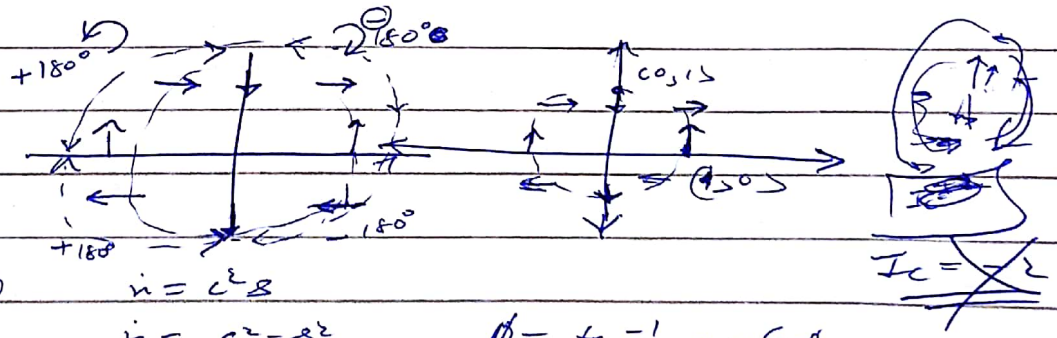
iii) saddle



iv) Center



$x = r \cos \theta$
 $y = r \sin \theta$
 $r = \sqrt{x^2 + y^2}$
 $\theta = \tan^{-1} \frac{y}{x}$



On unit circle

$x^2 + y^2 = 1$

- $(1, 0) \rightarrow x=1, y=0$
- $(0, 1) \rightarrow x=0, y=1$
- $(-1, 0) \rightarrow x=-1, y=0$
- $(0, -1) \rightarrow x=0, y=-1$

$I_c = 0$

Teacher's Sign: _____

$$x = 2n^2 - 1$$
$$y = 2ny$$

$$\tan \phi = \left(\frac{\sin 2\theta}{\cos 2\theta} \right) = \tan 2\theta$$
$$I_c = +2\phi$$

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